Granddad was superstitious about books. He thought that if you had enough of them around, education leaked out, like radioactivity.

Terry Pratchett

University of Notre Dame

JINA Lecture Series on Tools and Toys in Nuclear Astrophysics

Nuclear Reaction Network Techniques

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Los Alamos National Laboratory Steward Observatory, University of Arizona Syllabus

1 June 20 Purpose, Motivation, Forming a network, PP-chain code

2 June 21 Jacobian formation, Energy generation, Time integration, CNO-cycle code

3 June 22 Linear algebra, Thermodynamic trajectories, Alpha-chain code

4 June 23 Nuclear Statistical Equilibrium code, Big-Bang code

5 June 24 Networks in hydrodynamic simulations, General network code Sites of the week

nobelprize.org/physics/articles/fusion/index.html

ø hyperphysics.phy-astr.gsu.edu/hbase/astro/procyc.html

www.che.eng.ohio-state.edu/~FEINBERG/ LecturesOnReactionNetworks/

www.cococubed.com/papers/hix_meyer_2004.pdf

Purpose

The Joint Institute for Nuclear Astrophysics (JINA), a National Foundation for Science Physics Frontier Center, is organizing its second school on "Tools and Toys for Nuclear Astrophysics".

The school will focus on "Nuclear Reaction Network Techniques" and will take place at Notre Dame, IN between June 20 – July 1, 2005.

Reaction network techniques have been developed as an important tool for applications in nuclear astrophysics and many other physics areas.



Purpose

This technique is particularly important for simulating nucleosynthesis processes and for predicting element production as well as energy generation in both stellar evolution and stellar explosion processes.

The increasing need for using computational networking techniques requires special training in the development and application of network programs.

This school will provide an overview about the mathematical background as well as the computational network techniques.

Purpose

Different applications for stellar nucleosynthesis simulations will be presented and discussed.

In addition to the lectures, the participants will have time to work with the various network codes under the guidance of the lecturers.

Codes will be shared freely and will be available to participants after the workshop.

STELLAR INTERIORS SINS & NUCLEOSYNTHESIS

 $N_A = 6.022 \times 10^{23}$ number/mole = Avogadro's number

number = moles $\cdot N_A$ $N = M \cdot N_A$ $n = m \cdot N_A$

A_i = number of protons + neutrons in species i = atomic mass number = molar mass [g/mole]

 Z_i = number of protons in species i

- = charge
- = atomic number

 $mass = moles \cdot molar mass$

mass = $M \cdot A$ 1 mole of ¹²C has a mass of 12 g

 $\rho = m \cdot A$

Number density, mass density, mole fraction for a pure composition

$$n = \frac{\rho N_A}{A} \quad \text{cm}^{-3} \qquad \rho = \frac{n A}{N_A} \quad \text{g cm}^{-3}$$

 $Y = \frac{1}{A} = \frac{n}{\rho N_A}$ mole/g, molar fraction

For a mixture of isotopes,

 $Y_i = \frac{n_i}{\rho N_A}$ mole/g, molar fraction

 $X_i = A_i Y_i = \frac{A_i n_i}{\rho N_A} = \frac{\rho_i}{\rho}$ mass fraction, dimensionless

 $\sum X_i = 1$



 $\overline{A} = \frac{\sum n_i \overline{A_i}}{\sum n_i} = \frac{1}{\sum Y_i} \quad \text{g mole}^{-1} \qquad \overline{Z} = \frac{\sum n_i \overline{Z_i}}{\sum n_i} = \overline{A} \sum Y_i \overline{Z_i}$

$$n = \frac{\rho}{\overline{A}} N_A \quad \mathrm{cm}^{-3}$$

$$\rho = \frac{n}{N_A} \overline{A} \quad \text{g cm}^{-3}$$



Overview

These lectures will be about a special, but rather large class of ordinary differential equations (ODEs) – those that derive from nuclear/chemical/ionization/biological reaction networks.



the internet, sept3, 2004 9:22am UTC

In order to provide a sense of how these equations come about, it is useful to write down an example of reaction network and indicate informally how it induces a system of ODEs.

Α

A + C <

<---> 2B

This is a diagram for a reaction network. It indicates that species A can decompose into two units of species B, that two units of B can combine to form one unit of A, that a unit of species A can react with a unit of species C to form a unit of species D, and so on.

Δ

A + C <

 \leftarrow 2B

Suppose we throw the various species in a pot.

The pot is stirred constantly so that its contents remain spatially homogeneous for all time.

We'll also assume that the contents of the pot are forever maintained at constant temperature and total volume.



This is not to say the composition within the pot will remain constant in time, for the occurrence of reactions will consume some species and generate others. In fact, it is the temporal evolution of the composition that we wish to investigate.



Cosmic Cauldon Marilynn Flynn, 1995

Denote the instantaneous values of the molar abundances by Y_A, Y_B, Y_C, Y_D, and Y_E. We want to write down differential equations that describe the evolution of the five mole fractions.

The key to writing down the ODEs lies in knowing how rapidly each of the several reactions occurs. We assume the instantaneous rate of each reaction depends (in its own way) on the instantaneous composition vector Y.



• That is, we assume there exists non-negative, real valued rate function $\mathcal{K}_{A\to 2B}$ that gives the instantaneous rate of the reaction A -> 2B. Similarly we assume a rate function $\mathcal{K}_{2B\to A}$ for the reaction 2B -> A, a rate function $\mathcal{K}_{A+C\to D}$ for the reaction A+C -> D, and so on.

→ 2B

The kinetics for a reaction network is an assignment of a rate function for each reaction in the network.

The Let's begin by considering the instantaneous rate of change of Y_A . Every time the reaction A -> 2B occurs we lose a unit of species A, and that reaction has an occurrence rate of $\mathcal{K}_{A\rightarrow 2B}$

 ${\it o}$ On the other hand, every time 2B -> A occurs, we gain a unit of species A, and that reaction occurs at rate ${\cal K}_{2B\to A}$

 $A + C \iff D$

 $A \leftarrow 2B$

Similarly the reactions B+E -> A+C and D -> A+C produce a unit of species A with each occurrence, while the occurrence of A+C -> D results in a loss of a unit of species A. So we write

$$\dot{Y}_{A} = -\mathcal{K}_{A\to 2B} + \mathcal{K}_{2B\to A} - \mathcal{K}_{A+C\to D} + \mathcal{K}_{D\to A+C} + \mathcal{K}_{B+E\to A+C}$$



For species B, whenever A -> 2B occurs we gain two units of B, and whenever 2B -> A occurs we lose two units of B. When D -> B+E occurs we gain a unit of B, and we lose a unit of B when B+E -> A+C occurs. Thus,

 $Y_{\rm B} = 2\mathcal{K}_{\rm A\to 2B} - 2\mathcal{K}_{\rm 2B\to A} + \mathcal{K}_{\rm D\to B+E} - \mathcal{K}_{\rm B+E\to A+C}$



Continuing in this way we can write down equations for the other three species to generate the full system of ODEs that govern our reactor:

 $\dot{Y}_{A} = -\mathcal{K}_{A \to 2B} + \mathcal{K}_{2B \to A} - \mathcal{K}_{A+C \to D} + \mathcal{K}_{D \to A+C} + \mathcal{K}_{B+E \to A+C}$ $\dot{Y}_{\rm B} = 2\mathcal{K}_{\rm A\to 2B} - 2\mathcal{K}_{\rm 2B\to A} + \mathcal{K}_{\rm D\to B+E} - \mathcal{K}_{\rm B+E\to A+C}$ $\dot{Y}_{\rm C} = -\mathcal{K}_{\rm A+C\to D} + \mathcal{K}_{\rm D\to A+C} + \mathcal{K}_{\rm B+E\to A+C}$ $A \leftrightarrow 2B$ $A + C \iff D$ $\dot{Y}_{\rm D} = \mathcal{K}_{\rm A+C \rightarrow D} - \mathcal{K}_{\rm D \rightarrow A+C} - \mathcal{K}_{\rm D \rightarrow B+E}$ $Y_{\rm E} = \mathcal{K}_{\rm D \to B+E} - \mathcal{K}_{\rm B+E \to A+C}$

So far we haven't said anything about the nature of the rate functions. For nuclear reaction networks, the kinetics is taken to be of a mass action type.

For A -> 2B we assume the more A there is, the more occurrences of the reaction there will be. Thus, we take the instantaneous rate of A -> 2B to be proportional to Y_A:

 $\mathcal{K}_{\mathrm{A}\to 2\mathrm{B}} = \alpha Y_A$





Svante August Arrhenius Nobel Prize 1903

For A+C -> D, an occurrence requires a unit of species A meet a unit of species C in the reactor, and we take the probability of such an encounter to be proportional to the product Y_AY_C.

Although we don't presume that every encounter yields a unit of species D, we still take the occurrence rate to be given by

 $\mathcal{K}_{\mathrm{A+C}\to\mathrm{D}} = \gamma Y_A Y_C$



Skull cross section Leonardo da Vinci

With mass action kinetics, our rate functions take the form

 $\mathcal{K}_{A \to 2B} = \alpha Y_A$ $A \leftarrow \frac{\alpha}{\beta} \rightarrow 2B$ $\mathcal{K}_{2B\to A} = \beta Y_B^2$ $A + C \leftarrow \frac{\gamma}{\delta} \rightarrow D$ $\mathcal{K}_{A+C\to D} = \gamma Y_A Y_C$ 3 E $\mathcal{K}_{\mathrm{D}\to\mathrm{B}+\mathrm{E}} = \varepsilon Y_D$ $\mathcal{K}_{\mathrm{D}\to\mathrm{A+C}} = \delta Y_D$ B + E $\mathcal{K}_{\mathrm{B}+\mathrm{E}\to\mathrm{A}+\mathrm{C}} = \xi Y_B Y_E$

The rate constants may depend on temperature and density

 $A \leftarrow \frac{\alpha}{\beta} \rightarrow 2B$

3

 $A + C \leftarrow \frac{\gamma}{\delta} \rightarrow D$

E

And our reaction network takes the form

 $\dot{Y}_{A} = -\alpha Y_{A} + \beta Y_{B}^{2} - \gamma Y_{A}Y_{C} + \delta Y_{D} + \xi Y_{B}Y_{E}$

 $\dot{Y_{\rm B}} = 2\alpha Y_A - 2\beta Y_B^2 + \varepsilon Y_D - \xi Y_B \overline{Y_E}$

 $\dot{Y}_{\rm C} = -\gamma Y_A Y_C + \delta Y_D + \xi Y_B \overline{Y_E}$

 $\dot{Y}_{\rm D} = \gamma Y_A Y_C - \delta Y_D - \varepsilon Y_D$

 $\dot{Y}_{\rm E} = \varepsilon Y_D - \xi Y_B Y_E$

Interlude



Gerry Wasserburg 1993

Nuclear reaction networks

There are a number of types of nuclear reactions which are of astrophysical interest.

In addition to the emission or absorption of nuclei and nucleons, nuclear reactions can involve the emission or absorption of photons (γ-rays) and leptons (electrons, neutrinos, and their anti-particles).

As a result, nuclear reactions involve three of the four fundamental forces, the nuclear strong, electromagnetic and nuclear weak forces.

Weak interactions (those involving leptons) proceed much more slowly than those involving only nucleons and photons; however, these reactions are important because only weak interactions can change the global ratio of protons to neutrons.



The key quantity that describes the nuclear properties of a species important to nucleosynthesis or stellar energy generation is the cross section σ for interaction.

Consider the reaction i(j,k)l. The cross section σ_{ij} is the number of these reactions per target nucleus i per second divided by the flux of particles (number/cm²/s) of type j.

 $\sigma = \frac{\text{number of reactions target}^{-1} \text{sec}^{-1}}{\text{flux of incoming projectiles}} \quad \text{cm}^2$

• A representative unit is a barn, or 10^{-24} cm².

The reaction rate per unit volume, R_{ij} , in the simplest case is given by the flux of particles i times the number density of j times σ_{ij} .

 $R_{ij} = [\text{flux of } j] n_i \sigma_{ij} = v_{ij}n_j n_i \sigma_{ij} = n_i n_j \sigma_{ij} v_{ij}$

As the astrophysicist is concerned, nuclear physics boils down to the determination of the relevant cross sections (and the energy released by the reaction) which are usually averaged over a distribution of particle fluxes.

$$R_{i,j} = \int \sigma(|\vec{v}_i - \vec{v}_j|) |\vec{v}_i - \vec{v}_j| d^3 n_i d^3 n_j$$

We'll have more to say about the nature of cross sections and Q-values later on in this lecture series.

For now, we've established what we need to continue forming a nuclear reaction network.



OPPORTUNITIES IN NUCLEAR SCIENCE

A Long-Range Plan for the Next Decade

April 2002



The DOE/NSF Nuclear Science Advisory Committee U.S. Department of Energy • Office of Science • Division of Nuclear Physics National Science Foundation • Division of Physics • Nuclear Science Section

Let's start with a unidirectional binary reaction with unity coefficients.

 $i+j \rightarrow k+l$

 $\dot{Y}_i = -Y_i Y_j \overline{R_{ij}}$

 $\dot{Y}_j = -Y_i Y_j R_{ij}$

 $\dot{Y}_k = Y_i Y_j R_{ij}$

 $\dot{Y}_l = Y_i Y_j R_{ij}$

Now consider the case when the coefficients are not unity.

$$c_i i + c_j j \to c_k k + c_l l$$

$$\dot{Y}_i = -\frac{c_i}{c_i!c_j!}Y_i^{c_i}Y_j^{c_j}R_{ij}$$

$$\dot{Y}_j = -\frac{c_j}{c_i!c_j!}Y_i^{c_i}Y_j^{c_j}R_{ij}$$

$$\dot{Y}_k = \frac{c_k}{c_i!c_j!} Y_i^{c_i} Y_j^{c_j} R_{ij}$$

 $\dot{Y}_l = \frac{c_l}{c_i!c_j!} Y_i^{c_i} Y_j^{c_j} R_{ij}$

The factorial denominators avoid double counting. If identical reactants i=j (for example p + p) simply set c_j =0.

A general formula for a general binary reaction

 $c_i i + c_j j \leftrightarrow c_k k + c_l l$

$$\dot{Y}_{p} = \sum_{r,s} \frac{c_{p}}{c_{r}!c_{s}!} Y_{r}^{c_{r}} Y_{s}^{c_{s}} R_{rs} - \sum_{q} \frac{c_{p}}{c_{p}!c_{q}!} Y_{p}^{c_{p}} Y_{q}^{c_{q}} R_{pq}$$

If identical reactants i=j, simply set the $c_q=c_s=0$.

The reaction types can be divided into three categories based on the number of reactants which are nuclei.

Reactions involving a single nucleus, which include decays, electron and positron captures, photodisintegrations, and neutrino induced reactions, depend on the number density of only the target species.

$$\dot{Y}_i = \sum_j C_i R_j Y_j$$

The Ci's can be positive or negative numbers that specify how many particles of species i are created or destroyed.

For binary reactions, another general formula is

$$\dot{Y}_i = \sum_{jk} \frac{C_i}{C_j ! C_k !} \rho N_A R_{jk} Y_j Y_k$$



Binary sea exchange 2003

There are also a few important three-particle processes (like the triple-α process) which are commonly successive captures with an intermediate unstable target.

Using an equilibrium abundance for the unstable intermediate, the contributions of these reactions are commonly written in the form of a three-particle processes, depending on a trio of number densities.

$$\dot{Y}_{i} = \sum_{jkl} \frac{C_{i}}{C_{j}!C_{k}!C_{l}!} \rho^{2} N_{A}^{2} R_{jkl} Y_{j} Y_{k} Y_{l}$$

In terms of the reaction cross sections and molar abundances, a reaction network may described by the following set of ODEs

$$\dot{Y}_{i} = \sum_{j} C_{i} R_{j} Y_{j} + \sum_{jk} \frac{C_{i}}{C_{j}! C_{k}!} \rho N_{A} R_{jk} Y_{j} Y_{k} + \sum_{jkl} \frac{C_{i}}{C_{j}! C_{k}! C_{l}!} \rho^{2} N_{A}^{2} R_{jkl} Y_{j} Y_{k} Y_{l}$$



How does the Sun shine?

Wood – Ancient Greeks lasts 2000 years

Coal – Middle Ages lasts 4000 years

Gravitational – 1800's
 lasts 4 million years

 Nuclear reactions – 1940's lasts 10 billion years

How does the Sun shine?

Four hydrogen nuclei get transformed into one helium nucleus. The limiting step is a rare reaction; hence a long lived Sun.

But the mass of 4 hydrogen nuclei is larger than the mass of 1 helium nucleus. Where did the missing mass go?

 \oslash E = mc²

The sun presently shines by burning hydrogen (fuel) into helium (ash) in its core.



600 million tons of hydrogen are converted into 596 million tons of helium every second

Hans Bethe realized in 1939 that the weak interaction was capable of converting a proton into neutron during the brief encounter of a scattering event.





Since the neutron is more massive than a hydrogen nuclei, such a decay would require energy (endothermic) except for the fact that the neutron can appear in a bound state with proton in the form a deuterium nucleus.

The binding energy is sufficient (2.2245 MeV) to make the reaction exothermic.



We have four species to track (¹H, ²H, ³He, ⁴He), and three binary reactions that couple these species; p(p,e⁺ν)²H, ²H(p,γ)³He, and ³He(³He,2p)⁴He.



 $\dot{Y}_{\rm p} = -Y_{\rm p} \overline{Y_{\rm p} R_{\rm p,p}} - Y_{\rm p} \overline{Y_{\rm d} R_{\rm p,d}} + \overline{Y_{\rm 3he} Y_{\rm 3he} R_{\rm 3he,3he}}$

 $\dot{Y}_{\rm d} = 0.5Y_{\rm p}Y_{\rm p}R_{\rm p,p} - Y_{\rm p}Y_{\rm d}R_{\rm p,d}$

 $Y_{3he} = Y_p Y_d R_{p,d} - Y_{he3} Y_{he3} R_{he3,he3}$

 $\dot{Y}_{4he} = 0.5 Y_{he3} Y_{he3} R_{he3,he3}$



Where the density and Avogadro number dependence has been folded into the reaction rates

Prior to 1958 it was believed that the PPI chain would proceed under most conditions, even if an appreciable abundance of ⁴He were present.

Oblight Holmgren & Johnston measured the ³He(α,γ)⁷Be cross section to be 2500 times larger than the previously accepted value, so that this reaction competes with ³He(³He,2p)⁴He for ³He nuclei, particularly at elevated temperatures.

This possibility leads to two new chains for converting H to He, PPII and PPIII, and they correspond to the two possible fates of the ⁷Be nucleus.



The weights of the reactions are given for conditions in the Sun. The PP chains are the most important energy source in stars with masses less than 1.5 M_{sun} .

Tasks for the day

Derive the ODE equations for the PPI chain.

Download, compile, and run the pp-chain code from www.cococubed.com/code_pages/burn.shtml

Plot the abundances for temperatures of 1.0x10⁷ K, 2.0x10⁷ K, and 3.0x10⁷ K for a density of 150 g/cm³ and an initial composition of 75% H and 25% He by mass. Based on your calculations: How much hydrogen is left in the center of the Sun? How long will the Sun live? What do you conclude from the runs with different temperatures?

Tasks for the day

For those who enjoy a challenge, reproduce the following plot.



After Parker, Bahcall & Fowler ApJ 139, 602, 1964. Also Clayton figure 5–10.

Tools and Toys in Nuclear Astrophysics



Copper Cauldron with Three Eggs, 1734, oil on panel, Jean-Baptiste-Siméon Chardin