It is easier to believe that Yankee professors would lie than that stones would fall from heaven.

Thomas Jefferson

University of Notre Dame

JINA Lecture Series on Tools and Toys in Nuclear Astrophysics

Nuclear Reaction Network Techniques

Francis X. Timmes fxt44@mac.com cococubed.com/talk_pages/jina05.shtml

Los Alamos National Laboratory Steward Observatory, University of Arizona Sites of the week

theory.itp.ucsb.edu/~paxton/

Iash.uchicago.edu/website/home/

www.ucolick.org/~zingale/

Syllabus

1 June 20 Purpose, Motivation, Forming a network, PP-chain code

2 June 21 Jacobian formation, Energy generation, Time integration, CNO-cycle code

3 June 22 Linear algebra, Thermodynamic trajectories, Alpha-chain code

4 June 23 Nuclear Statistical Equilibrium code, Big-Bang code

5 June 24 Networks in hydrodynamic simulations, General network code

The constraints of mass and charge conservation



give two equations for the two unknowns, μ_n and μ_p .



If weak interactions are also balanced (e.g., neutrino capture occurring as frequently on the daughter nucleus as electron capture on the parent), then only two parameters, ρ and T, specify the abundances.



This last occurred for $T > 10^9$ K in the Big Bang.



 $x = \frac{m_e c^2}{kT}$

An ordinary differential equation for the photon temperature of the expanding universe.

 $g(x) = 1 + N_{\nu} \frac{7}{8} \left[\frac{4}{11} f(x) \right]^{4/3} + \int_{0}^{\infty} \sqrt{x^{2} + y^{2}} \left[\exp\left(\sqrt{x^{2} + y^{2}}\right) + 1 \right]^{-1} y^{2} dy$



How our universe cools down as it expands.



A typical Big Bang reaction network.



By 35 min nucleosynthesis is essentially complete.

Mass Fraction

The observed abundances of the light elements imply the density of normal matter in the universe is about 3.5 x 10⁻³¹ g/cm³.

 Four independent measurements of four different elements lead to a consistent constraint.

This gives us confidence that BBN provides a correct explanation of light element formation.

Baryon mass density

The Center for Astrophysical Thermonuclear Flashes, or Flash Center, was founded at the University of Chicago in 1997 under contract to the United States Department of Energy as part of its Accelerated Strategic Computing Initiative (ASCI).

The goal of the Center is to significantly advance the solution to several problems related to thermonuclear flashes on the surfaces and in the interiors of compact stars, in particular X-ray bursts, Type Ia supernovae, and classical novae.

Laser impnging on copper-foam target, Calder et al 2002

Helium burning on neutron stars, Zingale et al 2001

These problems are remarkable for the breadth of physical phenomena involved.

They range from accretion flow onto the surfaces of the compact stars to shear flow and Rayleigh-Taylor instabilities on the stellar surfaces, ignition of nuclear burning under conditions leading to convection, deflagration and detonation flame fronts, and stellar envelope expansion.

The physical processes include convection and turbulence at large Reynolds and Rayleigh numbers, convective penetration of stable matter at very high densities, equations of state for relativistic and degenerate matter, thermonuclear burning, mixing instabilities, burning front propagation, and radiation hydrodynamics.

Mixing at white dwarf surface, Alexakis et al,2004

Few astrophysical problems present a substantially greater level of physical complexity.

Cellular detonation Timmes et al 2001

sime = 150.453 ns number of blocks = 4736 AMR levels = 7

To advance the solution of these astrophysical problems requires the development of new simulation tools capable of handling the extreme resolution and physical requirements imposed by these thermonuclear flashes and to do so while making efficient use of the parallel supercomputers developed by the ASCI project, the most powerful constructed to date.

The FLASH code represents a step along the road to this goal.

Single mode RT instability, Calder et al. 2003

FLASH is a modular, adaptive, and parallel simulation code capable of handling general compressible flow problems in astrophysical environments.

It is modular: FLASH has been designed to allow users to configure initial and boundary conditions, change algorithms, and add new physical effects with minimal effort.

Off Center SNIa deflagration, Plewa et al. 2005

It is adaptive: FLASH uses a block-structured adaptive grid, placing resolution elements only where they are needed most.

		And the second se	-
-0.5 me = 0.000 µs umber of blocks =	0.0	0.5	1.0

Emery wind tunnel Ricker et al, 2000

It is parallel: FLASH uses the Message-Passing Interface library (mpich) to achieve portability and scalability on a variety of different message-passing parallel computers.

Jean's mass collapse Ricker et al, 2002, 2004

FLASH 2.5 Data Layout, Modularity, 2005

FLASH 0.0 Initial, 1999

FLASH 1.0 Organization, 2000

FLASH a hydrodynamics code for 1, 2 or 3D reactive flows.

Implementation in Fortran 90 + C + Pearl + IDL

~28,000 lines in driver + adaptive mesh refinement modules

~79,000 lines in physics modules

Object-oriented framework

Portable – runs efficiently on a wide class of machines

Parallel using Argonne's Message Passing Interface Library (mpich) library and uses the Hierarchical Data Format (HDF 5) or the Network Common Data Form (NETCDF) for parallel I/O

High-performance - Gordon Bell prize @ Supercomputing 2000

Ongoing verification and validation with analytic test problems Laboratory experiments

Available for distribution to the astrophysics community flash.uchicago.edu/website/home/

Reduces time to solution and improves accuracy by concentrating grid points in regions which require higher resolution

PARAMESH (NASA / GSFC) ct.gsfc.nasa.gov/paramesh/Users_manual/amr.html

 \odot Block structured refinement (8 x 8 x 8 blocks)

User-defined refinement criterion – currently using second derivatives of density and pressure

Local physics (hydro, eos, burn, etc.) occurs on a block. When guardcells are filled, operators act on each block as if isolated.

Number of guardcells depends on stencil size.

Number of interior points must balance tradeoff:
 More cells – more efficient (until block too big for cache)
 Fewer cells – can refine more quickly in smaller area.

Blocks and refinement are arranged in an 2^d-tree structure.

When a block is refined, 2^d child blocks are created, each with a factor of 2 resolution increase over its parent.

Blocks assigned indices via space-filling curve.

Neighboring blocks differ by at most one level of refinement.

Drawback: resolution can only fall of linearly in distance.

Feature: simplifies, speeds up accurate calculation of "boundary conditions" (guardcells)

Blocks are distributed
 across processors using a
 space-filling curve method.

A curve (typically a Morton curve) is threaded through the mesh, with weights according to how much each block "costs" in CPU use, communication, memory.

Other curves (Hilbert, Peano) with better clustering properties have been tried; no performance enhancement.

 Refinement and redistribution of blocks every four time steps.

zingale et al, 2000

The most common choice of the hydrodynamic module in the current version of the FLASH code solves Euler's equations for compressible gas dynamics. The equations can be written in conservative form as

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \tilde{\mathbf{v}} = 0$$

$$\frac{\partial \rho \tilde{\mathbf{v}}}{\partial t} + \nabla \cdot \rho \tilde{\mathbf{v}} \tilde{\mathbf{v}} + \nabla P = \rho \tilde{\mathbf{g}}$$

$$\frac{\partial \rho E}{\partial t} + \nabla \cdot \left(\rho E + P\right) \mathbf{\tilde{v}} = \rho \mathbf{\tilde{v}} \cdot \mathbf{\tilde{g}}$$

$$E = \epsilon + \frac{1}{2}v^2$$

The pressure is obtained from the internal energy and the density by an equation of state.

For the case of a simple gamma-law equation of state, the pressure is given by

 $P = (\gamma - 1)\rho\epsilon$

but more general equations of state are readily available.

Eos Goddess of Dawn Aurora to Romans

For reactive flows, a separate advection equation must be solved for each chemical or nuclear species

$$\frac{\partial \rho X_i}{\partial t} + \nabla \cdot \rho X_i \mathbf{\tilde{v}} = 0$$

The code does not explicitly track interfaces between the fluids, so that a small amount of numerical mixing can be expected during the course of the calculation.

Fryer et al, 2005

Physical and numerical errors are unavoidable in modeling thermonuclear reactive flows.

Physical errors are introduced by using approximations to known pieces of physics, or by using estimates for incompletely known physics.

An incompleteness error committed by the networks in FLASH are the approximations used for the reaction rates.

Many key rates aren't known experimentally to within a factor of 2-10.

How large an error is made in a stellar model that uses these imprecisely known reaction rate is highly problem dependent.

Another physical error made by the small reaction networks in FLASH is that they are all approximations to using a large and complete reaction network.

 As we've noted, a carefully designed α-chain and heavy ion reaction network can usually reproduce the energy generation rate of a large nuclear reaction network to within 30%.

For conditions where Y_e remains near 0.5 the accuracy is much better that 30%, but in regimes where Y_e differs from 0.5, the accuracy is significantly reduced.

Numerical errors are introduced by the algorithms and methods employed in a calculation.

The goal for the reaction networks in FLASH is to reduce the numerical integration error below the formal accuracy of the hydrodynamic algorithm; to have the accuracy of the hydrodynamics be the limiting source of numerical error.

Train wreck at Montparnasse, France, 1865

For fluid flows with very strong shock features, the absolute accuracy of the hydro algorithm is about two or three digits.

For smooth fluid flows the absolute accuracy is better, perhaps to five or six digits, but highly problem dependent.

The Weir Dhruva Mistry, 2005

The time integration routines in the reaction networks permit an integration accuracy that is well below the formal accuracy of the hydrodynamic algorithm, even for smooth fluid flows.

By calculating the approximate reactions networks as accurately as possible, we strive to eliminate the integration as a potential source of systematic numerical error.

FLASH is an explicit solver for compressible hydrodynamics. At best it is limited to time steps no larger than time it takes sound waves to cross a grid cell. So we are limited to "fast" phenomena – perhaps a few million sound crossing times.

There are many instances where one is interested on evolutions on much longer time scales, such as stars or various thermal timescale events in stars (e.g., slow neutron captures).

Richard Courant circa 1958

Bill Paxton

Stellar evolution code

Tioga

Kavli Institute for Theoretical Physics, UC Santa Barbara

Interlude

All the previous networks (big bang, hydrogen burners, alpha chains) are examples of hardwired networks.

Each of those networks are carefully crafted by hand. They have the advantage of being fast, but the disadvantage of being inflexible.

A general network, capable of doing any reaction network, is generally softwired.

Softwired 2005 kenn brown and chris wren

First, connect isotope i with isotope k by reaction type n

subroutine naray

c..this routine builds the nrr(7,i) array, which specifies c..the location of isotopes coupled to i by various reactions. c..the first index on nrr refers to reactions of the form c.. 1=ng 2=pn 3=pg 4=ap 5=an 6=ag 7=b-

c..initialize

jz(1) = 0 jz(2) = 1 jz(3) = 1 jz(4) = 1 jz(5) = 2 jz(6) = 2 jz(7) = -1 jn(1) = 1 jn(2) = -1 jn(3) = 0 jn(4) = 2 jn(5) = 1 jn(6) = 2jn(7) = 1

c..isotope i connected to isotope k by reaction type n

do i=ionbeg,ionend do n=1,7 nrr(n,i) = 0kz = int(zion(i)) + jz(n)kn = int(nion(i)) + jn(n)do k=ionbeg,ionend if (kz.eq.int(zion(k)) .and. kn.eq.int(nion(k))) nrr(n,i)=k enddo enddo enddo

return end

Second, form the ODEs.

c..for every isotope in the network do j=ionbeg,ionend

c..set up the y(j)(n,g)y(k) and y(k)(g,n)y(j) components k = nrr(1,j)if (k.gt. 0) then $b1 = -aan^*sig(1,j)^*y(j) + sig(2,j)^*y(k)$ dydt(j) = dydt(j) + b1 dydt(in) = dydt(in) + b1 dydt(k) = dydt(k) - b1end if

Third, form the Jacobian

Forth, evolve the system in time like any of our hardwired networks. c..for every isotope in the network do j=ionbeg,ionend

c...set up the (n,g) components $\mathbf{k} = \operatorname{nrr}(1,\mathbf{j})$ if (k . gt. 0) then a1 = sig(1,j) * aana2 = sig(2,j)a3 = sig(1,j) * y(j)dfdy(j,j) = dfdy(j,j) - a1dfdy(j,k) = dfdy(j,k) + a2dfdy(j,in) = dfdy(j,in) - a3dfdy(k,j) = dfdy(k,j) + a1dfdy(k,k) = dfdy(k,k) - a2dfdy(k,in) = dfdy(k,in) + a3dfdy(in,j) = dfdy(in,j) - a1dfdy(in,k) = dfdy(in,k) + a2dfdy(in,in) = dfdy(in,in) - a3end if

Tasks for the day

Download, compile, and run the torch code from www.cococubed.com/code_pages/burn.shtml How do you set which isotopes are in the network?

Run one of the task problems from the previous four days with the torch network. Do you get the same answers? Why or why not?

Extra credit: The reading the nuclear data file is a bit clumsy. Modify the code to use the NON-SMOKER data base of reaction rates. Modify the code to use the LAnganke et al weak reaction rates.

Tools and Toys in Nuclear Astrophysics

